

LT. Laplace Transform

1. Translation formula. The usual L.T. formula for translation on the t -axis is

$$(1) \quad \mathcal{L}(u(t-a)f(t-a)) = e^{-as}F(s), \quad \text{where } F(s) = \mathcal{L}(f(t)), \quad a > 0.$$

This formula is useful for computing the inverse Laplace transform of $e^{-as}F(s)$, for example. On the other hand, as written above it is not immediately applicable to computing the L.T. of functions having the form $u(t-a)f(t)$. For this you should use instead this form of (1):

$$(2) \quad \mathcal{L}(u(t-a)f(t)) = e^{-as}\mathcal{L}(f(t+a)), \quad a > 0.$$

Example 1. Calculate the Laplace transform of $u(t-1)(t^2+2t)$.

Solution. Here $f(t) = t^2 + 2t$, so (check this!) $f(t+1) = t^2 + 4t + 3$. So by (2),
$$\mathcal{L}(u(t-1)(t^2+2t)) = e^{-s}\mathcal{L}(t^2+4t+3) = e^{-s}\left(\frac{2}{s^3} + \frac{4}{s^2} + \frac{3}{s}\right).$$

Example 2. Find $\mathcal{L}(u(t-\frac{\pi}{2})\sin t)$.

Solution.
$$\begin{aligned} \mathcal{L}(u(t-\frac{\pi}{2})\sin t) &= e^{-\pi s/2}\mathcal{L}(\sin(t+\frac{\pi}{2})) \\ &= e^{-\pi s/2}\mathcal{L}(\cos t) = e^{-\pi s/2}\frac{s}{s^2+1}. \end{aligned}$$

Proof of formula (2). According to (1), for any $g(t)$ we have

$$\mathcal{L}(u(t-a)g(t-a)) = e^{-as}\mathcal{L}(g(t));$$

this says that to get the factor on the right side involving g , we should replace $t-a$ by t in the function $g(t-a)$ on the left, and then take its Laplace transform.

Apply this procedure to the function $f(t)$, written in the form $f(t) = f((t-a)+a)$; we get (“replacing $t-a$ by t and then taking the Laplace Transform”)

$$\mathcal{L}(u(t-a)f((t-a)+a)) = e^{-as}\mathcal{L}(f(t+a)),$$

exactly the formula (2) that we wanted to prove. □

Exercises. Find: a) $\mathcal{L}(u(t-a)e^t)$ b) $\mathcal{L}(u(t-\pi)\cos t)$ c) $\mathcal{L}(u(t-2)te^{-t})$

Solutions. a) $e^{-as}\frac{e^a}{s-1}$ b) $-e^{-\pi s}\frac{s}{s^2+1}$ c) $e^{-2s}\frac{e^{-2}(2s+3)}{(s+1)^2}$

2. The transfer function and Green's function.

If we use the Laplace transform to solve the IVP

$$y'' + ay' + by = r(t), \quad y(0) = 0, \quad y'(0) = 0,$$

the transform of the IVP, with the usual notation, is

$$s^2Y + asY + bY = R(s);$$

whose solution for $Y = \mathcal{L}^{-1}(y)$ is

$$Y = R(s) \frac{1}{s^2 + as + b};$$

using the convolution operator to take the inverse transform, we get as the solution (further down the function $w(t)$ is defined):

$$(3) \quad y = r(t) * w(t) = \int_0^t r(u)w(t-u) du .$$

In this form of the solution, the following terminology is often used. Let $p(D) = D^2 + aD + b$ be the differential operator; then we write

$$\begin{aligned} W(s) &= \frac{1}{s^2 + as + b} && \text{the **transfer function** for } p(D), \\ w(t) &= \mathcal{L}^{-1}(W(s)) && \text{the **weight function** for } p(D), \\ G(t, u) &= w(t-u) && \text{the **Green's function** for } p(D). \end{aligned}$$

The important thing to note is that each of these depends only on the operator, not on the forcing function $r(t)$; once they are calculated, the solution (3) to the IVP can be written down immediately as the definite integral there, and used for a variety of different $r(t)$.

The weight function $w(t)$ can be thought of as the unique solution to the IVP

$$(4) \quad y'' + ay' + by = 0; \quad y(0) = 0, \quad y'(0) = 1;$$

or as the solution to the IVP

$$(5) \quad y'' + ay' + by = \delta(t); \quad y(0) = 0, \quad y'(0^-) = 0;$$

in the second equation, $\delta(t)$ is the Dirac delta function. Both IVP's model (for $a, b > 0$) a damped spring-mass system which is initially at rest, but whose mass is given a unit impulse at time zero, say by a sharp blow.

It is an easy exercise to show that $w(t)$ is the solution to both IVP's. As an example of Green's functions, see the last few Laplace Transform exercises (in Section 3D).

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